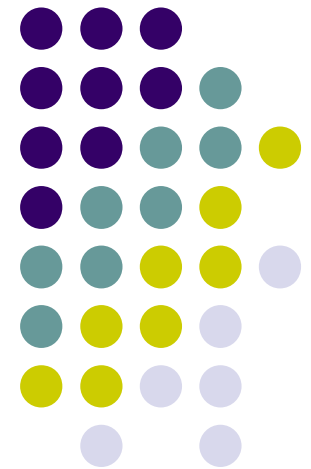
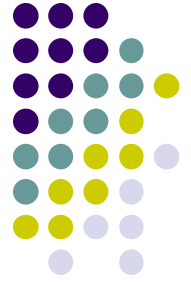


$$V \Leftrightarrow \vec{E}$$

1D example

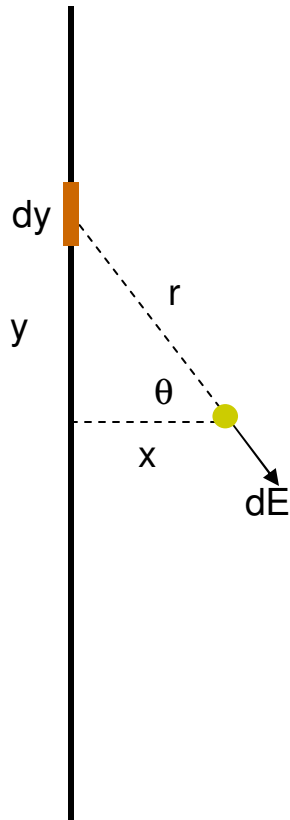
Dr. Ray Kwok  
SJSU





# 1D uniform charge distribution

At center line of a finite length  $L$ ,  $E = ?$



$$|dE| = k \frac{|dq|}{r^2} = k \frac{(\lambda dy)}{x^2 + y^2}$$

$$|dE|_x = |dE| \cos \theta = k \frac{(\lambda dy)}{x^2 + y^2} \left( \frac{x}{r} \right) = \frac{k \lambda x (dy)}{(x^2 + y^2)^{3/2}}$$

$$E = \int_{-L/2}^{L/2} \frac{k \lambda x (dy)}{(x^2 + y^2)^{3/2}}$$

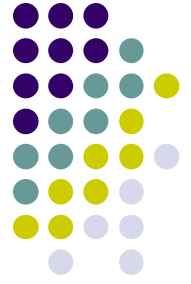
$$\int \frac{dy}{r^3} = \int \frac{x \sec^2 \theta d\theta}{(x \sec \theta)^3} = \frac{1}{x^2} \int \cos \theta d\theta = \frac{1}{x^2} \sin \theta = \frac{1}{x^2} \left[ \frac{y}{\sqrt{x^2 + y^2}} \right]$$

$$\int_{-L/2}^{L/2} \frac{dy}{r^3} = \frac{1}{x^2} \left[ \frac{y}{\sqrt{x^2 + y^2}} \right]_{-L/2}^{L/2} = \frac{1}{x^2} \frac{L}{\sqrt{x^2 + (L/2)^2}} = \frac{1}{x^2} \frac{L}{(L/2) \sqrt{(2x/L)^2 + 1}} = \frac{2}{x^2} \left( 1 + \left( \frac{2x}{L} \right)^2 \right)^{-1/2}$$

$$E = k \lambda x \cdot \frac{2}{x^2} \left( 1 + \left( \frac{2x}{L} \right)^2 \right)^{-1/2}$$

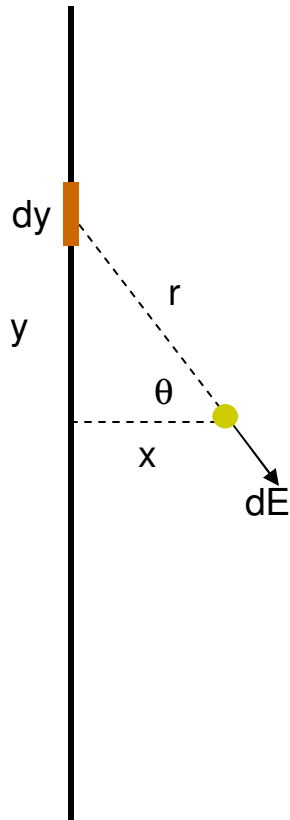
$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 x} \left( 1 + \left( \frac{2x}{L} \right)^2 \right)^{-1/2} \hat{x} = \frac{\lambda}{2\pi\epsilon_0 x \sqrt{1 + (2x/L)^2}} \hat{x}$$

(eqn. 1)



# Infinite line charge

Infinite length.  $E = ?$



From Eqn. 1, set  $L \rightarrow \infty$

$$\vec{E} = \lim_{L \rightarrow \infty} \frac{\lambda}{2\pi\epsilon_0 x} \left( 1 + \left( \frac{2x}{L} \right)^2 \right)^{-1/2} \hat{x}$$

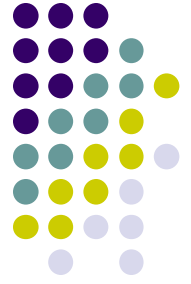
$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 x} \hat{x} \quad (\text{eqn. 2})$$

which can be easily confirmed with Gauss Law

$$EA = \frac{Q_{\text{inside}}}{\epsilon_0}$$

$$E(2\pi rL) = \frac{\lambda L}{\epsilon_0}$$

$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 r} \hat{r}$$



# point charge

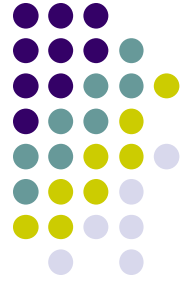
$L \rightarrow 0, E = ?$

From Eqn. 1, set  $L \rightarrow 0$

$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 x \sqrt{1 + (2x/L)^2}} \hat{x} = \frac{\lambda L}{2\pi\epsilon_0 x 2x \sqrt{(L/2x)^2 + 1}} \hat{x}$$

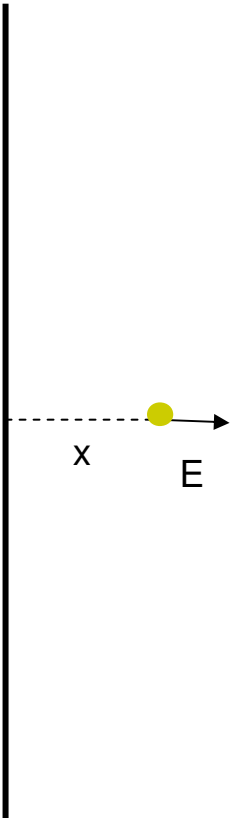
$$\vec{E} = \frac{\lambda L}{4\pi\epsilon_0 x^2} \hat{x} = \frac{Q}{4\pi\epsilon_0 x^2} \hat{x}$$

which is just the point charge equation.



# V from E

due to the charges on finite length L. (note:  $V = 0$  at infinity.)



$$\int_0^V dV = -\int_{\infty}^x \vec{E} \cdot d\vec{\ell}$$

$$V = -\int_{\infty}^x \frac{\lambda}{2\pi\epsilon_0 x} \left( 1 + \left( \frac{2x}{L} \right)^2 \right)^{-1/2} dx$$

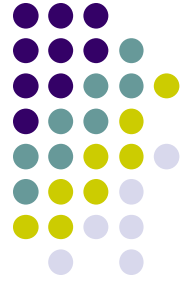
$$u \equiv \frac{2x}{L}$$

$$V = -\frac{\lambda}{2\pi\epsilon_0} \int_{\infty}^u \frac{du}{u\sqrt{1+u^2}}$$

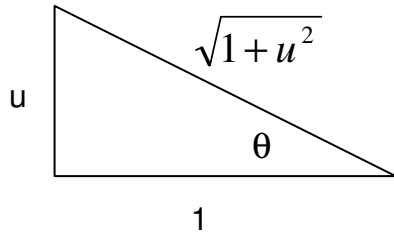
(see next page)

$$V = -\frac{\lambda}{2\pi\epsilon_0} \ln \left| \frac{u}{1 + \sqrt{1+u^2}} \right| = \frac{\lambda}{2\pi\epsilon_0} \ln \left| \frac{1 + \sqrt{1 + (2x/L)^2}}{2x/L} \right|$$

(eqn.3)



# integral



$$\int_{\infty}^u \frac{du}{u\sqrt{1+u^2}} = ?$$

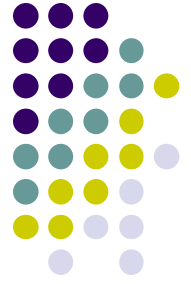
$$u = \tan \theta$$

$$du = \sec^2 \theta d\theta$$

$$\sqrt{1+u^2} = \sec \theta$$

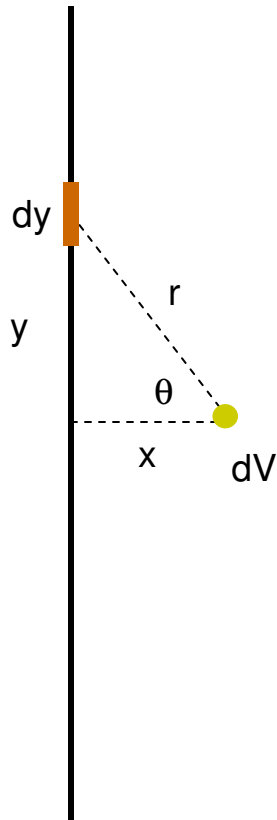
$$\int \frac{du}{u\sqrt{1+u^2}} = \int \frac{\sec^2 \theta d\theta}{\tan \theta \sec \theta} = \int \frac{d\theta}{\sin \theta} = \int \csc \theta d\theta = -\ln|\csc \theta + \cot \theta|$$

$$\int_{\infty}^u \frac{du}{u\sqrt{1+u^2}} = -\ln \left| \frac{\sqrt{1+u^2}}{u} + \frac{1}{u} \right|_{\infty}^u = \ln \left| \frac{u}{1+\sqrt{1+u^2}} \right|_{\infty}^u = \ln \left| \frac{u}{1+\sqrt{1+u^2}} \right|$$



# V from integrating pt charges

Finite length  $L$ .  $V = ?$



$$dV = k \frac{|dq|}{r} = k \frac{(\lambda dy)}{\sqrt{x^2 + y^2}}$$

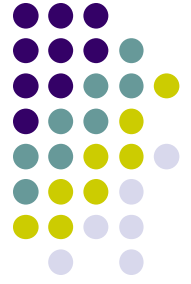
$$V = \int_{-L/2}^{L/2} \frac{k\lambda dy}{\sqrt{x^2 + y^2}}$$

$$\int \frac{dy}{r} = \int \frac{x \sec^2 \theta d\theta}{(x \sec \theta)} = \int \sec \theta d\theta = \ln |\sec \theta + \tan \theta| = \ln \left| \frac{r}{x} + \frac{y}{x} \right|$$

$$\int_{-L/2}^{L/2} \frac{dy}{r} = \ln \left| \frac{\sqrt{x^2 + y^2} + y}{x} \right|_{-L/2}^{L/2} = \ln \left| \frac{\sqrt{x^2 + (L/2)^2} + L/2}{\sqrt{x^2 + (L/2)^2} - L/2} \right| = \ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right|$$

$$V = \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right| \quad (\text{eqn. 4})$$

Note: Equations 3 & 4 are essentially the same !!!!!



# E from V – finite linear charge

From eqn (3)

$$V = \frac{\lambda}{2\pi\epsilon_0} \ln \left| \frac{1 + \sqrt{1 + (2x/L)^2}}{2x/L} \right| = \frac{\lambda}{2\pi\epsilon_0} \left[ \ln |1 + \sqrt{1 + u^2}| - \ln |u| \right]$$

$$u = 2x/L$$

$$E = -\frac{dV}{dx} = -\frac{2}{L} \frac{dV}{du} = -\frac{\lambda}{\pi\epsilon_0 L} \left[ \frac{\frac{1}{2}(1+u^2)^{-1/2}(2u)}{1 + \sqrt{1+u^2}} - \frac{1}{u} \right] = -\frac{\lambda}{\pi\epsilon_0 L} \left[ \frac{u^2 - \sqrt{1+u^2}(1 + \sqrt{1+u^2})}{u\sqrt{1+u^2}(1 + \sqrt{1+u^2})} \right]$$

$$\vec{E} = -\hat{x} \frac{\lambda}{\pi\epsilon_0 L} \left[ \frac{u^2 - \sqrt{1+u^2} - 1 - u^2}{u\sqrt{1+u^2}(1 + \sqrt{1+u^2})} \right] = \hat{x} \frac{\lambda}{\pi\epsilon_0 L u \sqrt{1+u^2}} = \frac{\lambda}{2\pi\epsilon_0 x \sqrt{1 + (2x/L)^2}} \hat{x}$$

From eqn (4)

$$V = \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right| = \frac{\lambda}{4\pi\epsilon_0} \left[ \ln |\sqrt{u^2 + 1} + 1| - \ln |\sqrt{u^2 + 1} - 1| \right]$$

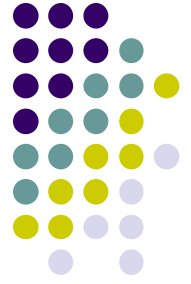
$$\vec{E} = -\hat{x} \frac{dV}{dx} = -\hat{x} \frac{2}{L} \frac{dV}{du} = -\hat{x} \frac{\lambda}{2\pi\epsilon_0 L} \left[ \frac{\frac{1}{2}(u^2 + 1)^{-1/2}(2u)}{\sqrt{u^2 + 1} + 1} - \frac{\frac{1}{2}(u^2 + 1)^{-1/2}(2u)}{\sqrt{u^2 + 1} - 1} \right]$$

$$\vec{E} = -\hat{x} \frac{\lambda}{2\pi\epsilon_0 L} \frac{u}{\sqrt{u^2 + 1}} \left[ \frac{(\sqrt{u^2 + 1} - 1) - (\sqrt{u^2 + 1} + 1)}{u^2 + 1 - 1} \right] = \hat{x} \frac{2\lambda}{2\pi\epsilon_0 L u \sqrt{u^2 + 1}}$$

$$\vec{E} = \hat{x} \frac{\lambda}{2\pi\epsilon_0 x \sqrt{1 + (2x/L)^2}}$$

Both give the same E as in eqn (1)



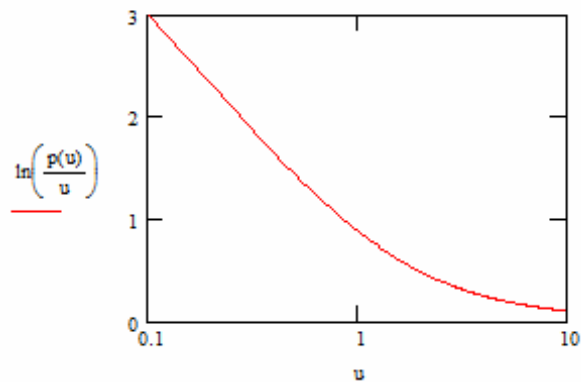
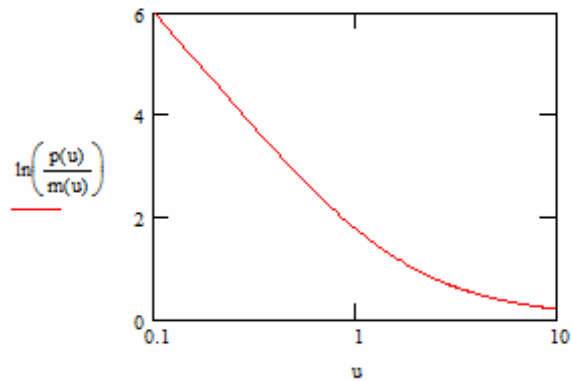


# Eqn(3) is the same as Eqn(4)

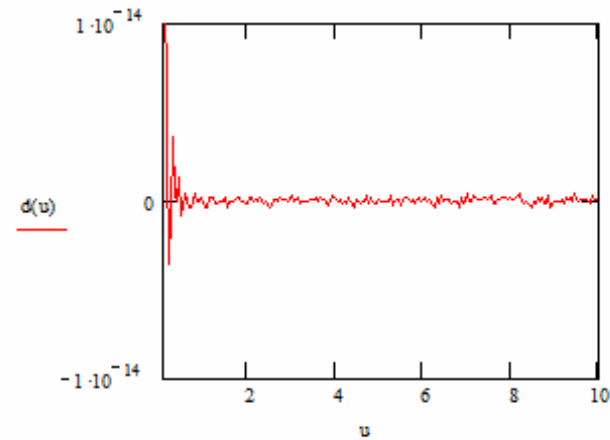
$u := 0.1, 0.15 \dots 10$

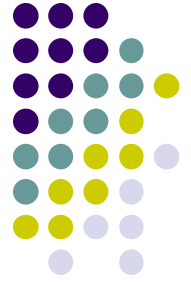
$$p(u) := \sqrt{1 + u^2} + 1$$

$$m(u) := \sqrt{1 + u^2} - 1$$



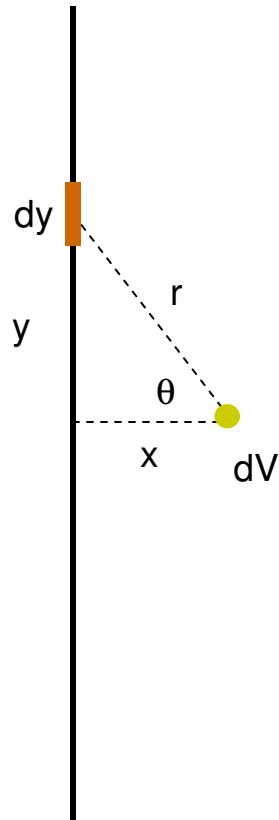
$$d(u) := \ln\left(\frac{p(u)}{m(u)}\right) - 2 \cdot \ln\left(\frac{p(u)}{u}\right)$$





# V from infinite line charge

From eqn (4)



$$V = \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right|$$

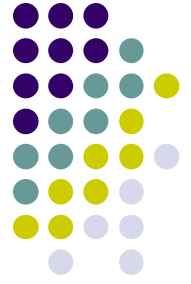
$$L \rightarrow \infty$$

$$\sqrt{(2x/L)^2 + 1} = \left(1 + (2x/L)^2\right)^{1/2} \approx 1 + 2\left(\frac{x}{L}\right)^2$$

$$\ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right| \approx \ln \left| \frac{2}{2(x/L)^2} \right| = 2 \ln \left| \frac{L}{x} \right| = -2 \ln |x| + C$$

$$V = -\frac{\lambda}{2\pi\epsilon_0} \ln |x| + C \quad \text{(eqn. 5)}$$

even though the constant  $\rightarrow \infty$



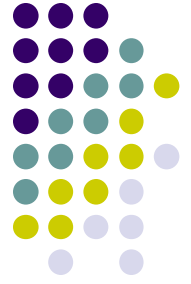
# V from E: Infinite line charge

From eqn (3)

$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 x} \hat{x}$$

$$V = -\int \vec{E} \cdot d\vec{\ell} = -\frac{\lambda}{2\pi\epsilon_0} \int \frac{dx}{x} = -\frac{\lambda}{2\pi\epsilon_0} \ln|x| + C$$

which is the same as eqn (5)



# V of a point charge

From eqn (4), set  $L \rightarrow 0$

$$V = \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{\sqrt{(2x/L)^2 + 1} + 1}{\sqrt{(2x/L)^2 + 1} - 1} \right| = \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{\sqrt{1 + (L/2x)^2} + (L/2x)}{\sqrt{1 + (L/2x)^2} - (L/2x)} \right| \approx \frac{\lambda}{4\pi\epsilon_0} \ln \left| \frac{1 + (L/2x)}{1 - (L/2x)} \right|$$

$$\ln |1 + x| = x - \frac{x^2}{2} + \frac{x^3}{3} - \frac{x^4}{4} + \dots \approx x$$

$$\ln \left| \frac{1 + (L/2x)}{1 - (L/2x)} \right| = \ln \left| 1 + \frac{2(L/2x)}{1 - (L/2x)} \right| \approx \frac{2(L/2x)}{1 - (L/2x)} = \frac{2L}{2x - L} \approx \frac{L}{x}$$

$$V = \frac{\lambda}{4\pi\epsilon_0} \left( \frac{L}{x} \right) = \frac{Q}{4\pi\epsilon_0 x}$$

which is just the point charge equation.