## MUON SPIN DEPOLARIZATION IN Gd- AND EuBa<sub>2</sub>Cu<sub>3</sub>Ox

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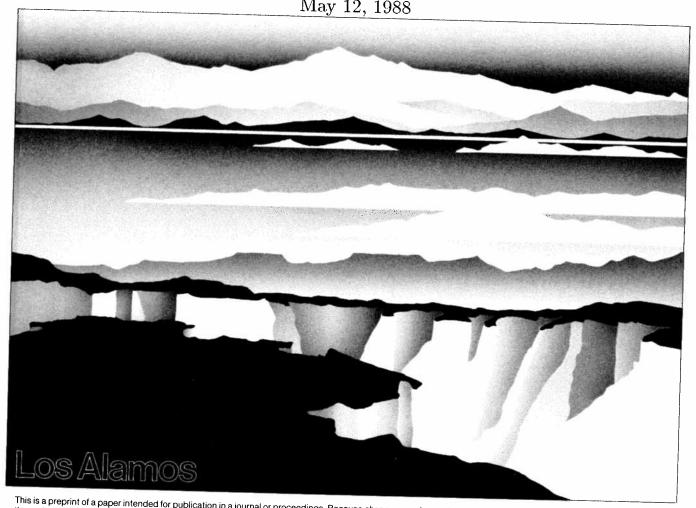
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> > LA-UR-88-1519

May 12, 1988



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SUBMITTED TO:

Proceedings of the International Conference on the First Two Years of High-Temperature Superconductivity Tuscaloosa, AL

April 11-13, 1988

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#### MUON SPIN DEPOLARIZATION IN Gd- and EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>

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Abstract Positive muon spin rotation ( $\mu$ SR) measurements on Gd- and EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> ( $x\approx7$ ) have been conducted in the temperature interval 4 - 300K. For each sample, muons stop both at grain boundaries and within the superconducting grains. Measured magnetic field penetration depths are 1550 and 1900Å for two specimens of GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>, and 1350Å for EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>.

Standard  $\mu$ SR techniques¹ were used to obtain positive muon ( $\mu^+$ ) Gaussian depolarization rates  $\Lambda$  (T) for Eu- and GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> in a 1 kOe transverse field. Also, zero-field data were taken for the latter sample. Figures 1 and 2 show  $\Lambda$ (T) and the corresponding muon precessional frequencies  $\nu$ (T) for a well-characterized polycrystalline sample of GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>. Two relaxation rates are observed for T<T<sub>c</sub>, corresponding to two distinct muon stopping regions. From the  $\nu$ (T) data we conclude that muons are thermalized in both normal and superconducting volumes of the sample. Further, the measured Fourier power and asymmetry data indicate that the superconducting to normal volume ratio is 5.6. We suggest that

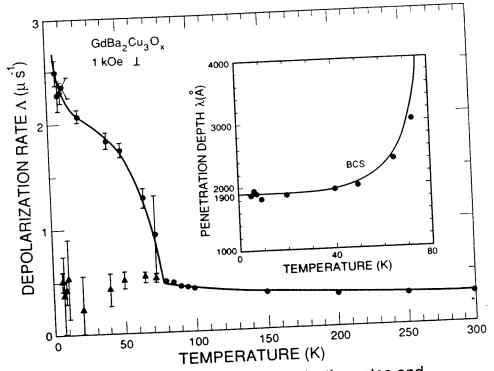


FIGURE 1. Muon Gaussian depolarization rates and penetration depth in GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>.

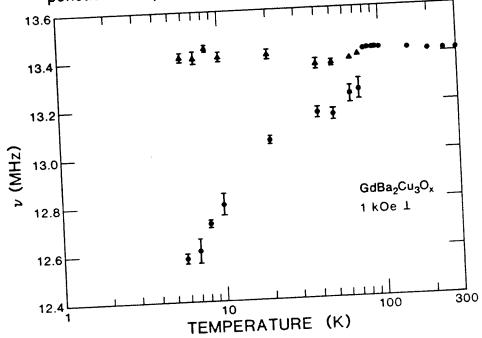


FIGURE 2. Muon frequencies in  $GdBa_2Cu_3O_x$ .

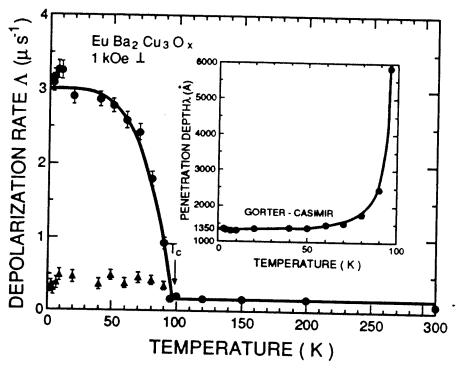


FIGURE 3. Muon Gaussian depolarization rates and penetration depth in EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>.

85% of stopped muons reside within the superconducting grains and that 15% occupy grain boundaries. Similar results are obtained from polycrystalline  $EuBa_2Cu_3O_x$  as shown in Figs. 3 and 4, and also from a second sample of  $GdBa_2Cu_3O_x$  (not shown).

The magnetic field penetration depth  $\lambda(T)$  is related to the field inhomogeneity, and thus  $\Lambda(T)$ , by the expression

$$< |\Delta H|^2 > = (B\phi/4\pi\lambda^2)[1+(4\pi^2\lambda^2B)/\phi]^{-1} = 2\Lambda^2/\gamma_{\mu}^2$$
 (1)

where  $\Lambda$  is the  $\mu^+$  depolarization rate,  $\gamma_\mu$  is the muon gyromagnetic ratio,  $\lambda$  is magnetic penetration depth, and  $\varphi$  is the flux quantum. By properly extrapolating the measured depolarization rates of Figs. 1 and 3 to zero Kelvin, and



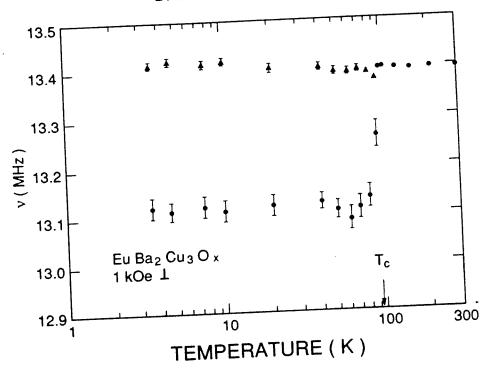


FIGURE 4. Muon frequencies in EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>.

applying eq. (1), we find  $\lambda(0)$  to be 1550 and 1900Å for two samples of GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>, and 1350Å for EuBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub>. Moreover, we compute the temperature dependence of  $\lambda(T)$  by assuming a phenomenological <sup>2</sup> expression  $\lambda(T) = \lambda(0)[1 - (T/T_c)^4]^{-1/2}$  which we substitute into eq. (1), with the above measured values for  $\lambda(0)$ . These results are shown in the insets of Figs. 1 and 3. The solid lines represent the phenomenological expression calculated with the appropriate  $T_c$ 's.

Zero-field  $\mu^+$  exponential relaxation rates and asymmetries for  $GdBa_2Cu_3O_{\chi}$  are shown in Fig. 5. With decreasing temperature there is a drop in asymmetry and concomitant increase in  $\Lambda$  near  $T_C$ . Enhanced field inhomogeneity can explain the increase in  $\Lambda,$  but not the

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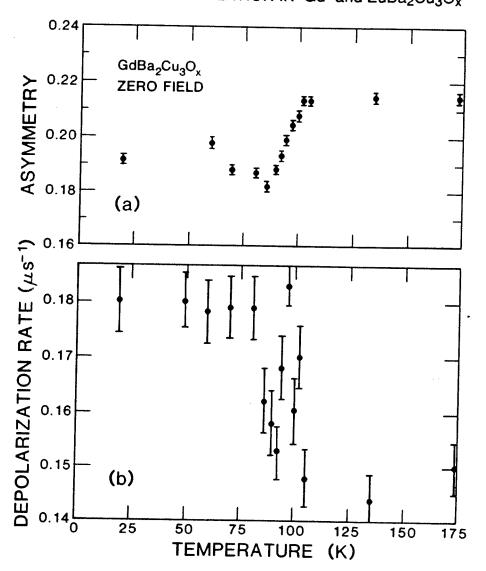


FIGURE 5. Muon exponential depolarization rates and asymmetries in GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> taken in zero field.

decrease in asymmetry; thus we offer no explanation for this puzzling result.

Work at Los Alamos is conducted under the aegis of the United States Department of Energy. Research support by Associated Western Universities, San Jose State Foundation, and the Robert A. Welch Foundation is acknowledged.

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